

Fundamentals of QM

- Physical system
- Quantum states
- Unitary evolution
- Measurement

A quantum (physical) system corresponds to a Hilbert space $\mathcal{H}$.

$\mathcal{H}$ infinite-dimensional —
"unbound" systems, e.g.
unconfined particles, fields, etc.

$\mathcal{H}$ finite-dimensional ($\mathbb{C}$ -space):
bound systems — finite
degrees of freedom:
confined particles
angular momentum
— orbital and momentum
in an atom
— intrinsic spin

Fix $\mathcal{H}$ $\mathbb{C}$ -space for a system.

At any given time the system is in a state represented by a unit vector $u \in \mathcal{H}$.

$$\|u\|=1.$$

Digression: Dirac Notation (column)

A unit vector (state) is written as $|\psi\rangle$

(a ket). ψ is some label.

Its adjoint $|\psi\rangle^* = \langle\psi|$

(a bra). If $|\psi\rangle$ and

$|\phi\rangle$ are kets we can form the inner product

$$\langle\phi|\psi\rangle \in \mathbb{C}$$

("bracket")

$|\psi\rangle\langle\phi|$ is an operator (if $\mathcal{H}$ is n-dimensional, this is an $n \times n$ matrix) 

Unitary evolution:

An isolated system evolves according to a (time-dependent) unitary operator:

$$U(t_2, t_1)$$

$t_1, t_2 \in \mathbb{R}$ moments in time and $U(t_2, t_1)$ is unitary. So that, if the system is in state $|\psi\rangle$ at time t_1 , it will be in state

$$U(t_2, t_1)|\psi\rangle$$

at time t_2 . (not necessary that $t_2 > t_1$).

$$\| |\psi\rangle \| = 1 \Rightarrow \| U(\cdot) |\psi\rangle \| = 1$$

b/c $U(\dots)$ is unitary. So states get mapped to states. Additional properties (common)

$$U(t, t) = I \quad (\forall t)$$

$$U(t_2, t_1) = U(t_1, t_2)^* \quad \left[= U(t_1, t_2)^{-1} \right]$$

$$\forall t_1, t_2$$

$$\forall t_1, t_2, t_3$$

$$U(t_3, t_2)U(t_2, t_1) = U(t_3, t_1)$$

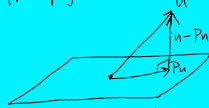
Measurement: Projective measurements

Def:
An operator $P \in \mathcal{L}(\mathcal{H})$ is a projector (orthogonal projection operator) if

$$1) P = P^* \quad (P \text{ is Hermitian})$$

$$2) P = P^2 \quad (P \text{ is idempotent})$$

$\mathcal{H}$: P projects $\mathcal{H}$ onto a subspace $\mathcal{U}$



P projects orthogonally:
show that $\langle u - Pu, Pu \rangle = 0$.

$$\begin{aligned} \langle u - Pu, Pu \rangle &= \langle u, Pu \rangle - \langle Pu, Pu \rangle \\ &= \langle u, Pu \rangle - \langle u, P^* Pu \rangle \\ &= \langle u, Pu \rangle - \langle u, PPu \rangle \\ &= \langle u, Pu \rangle - \langle u, Pu \rangle = 0 \end{aligned}$$

Note: I is a projector
 O " " "

Def: A complete set of orthogonal projectors (csop)

is a set $\{P_1, \dots, P_k\} \subseteq \mathcal{L}(\mathcal{H})$ such that
- each P_i is a nonzero projector
- $\sum_{i=1}^k P_i = I$

Ex: If $\{P_i\}$ is a csop then $P_i P_j = P_j P_i = 0$ for all $i \neq j$.

Def: Let $\mathcal{H}$ be the $\mathbb{C}$ -space of a physical system.

A projective measurement of the system is modeled by a csop $\{P_1, \dots, P_k\} \subseteq \mathcal{L}(\mathcal{H})$

Suppose the system is in some state $|\psi\rangle$ then measuring the system yields one of the outcomes $1, \dots, k$ with probability

$$Pr[j] = \|P_j |\psi\rangle\|^2$$

Projectors correspond to the possible outcomes of the measurement (a random experiment with sample space $\{1, \dots, k\}$).

Notice that $\forall j, 1 \leq j \leq k$

$$\begin{aligned} Pr[j] &= \|P_j |\psi\rangle\|^2 \\ &= \langle P_j |\psi\rangle, P_j |\psi\rangle \rangle \end{aligned}$$

Supposing the outcome is some j ($1 \leq j \leq k$), the post-measurement state is

$$\frac{P_j |\psi\rangle}{\|P_j |\psi\rangle\|}$$

$$= \frac{P_j |\psi\rangle}{\sqrt{P_j}}$$

Note $\left\| \frac{P_j |\psi\rangle}{\|P_j |\psi\rangle\|} \right\|$

$$= \|P_j |\psi\rangle\|^{-1} \|P_j |\psi\rangle\| = 1$$

Check: $P_r[j] \geq 0$ because it is a squared norm.

Need: $\sum_{j=1}^k P_r[j] = 1$:

$$\begin{aligned} \sum_{j=1}^k P_r[j] &= \sum_{j=1}^k \langle \psi | P_j | \psi \rangle \\ &= \langle \psi | \left(\sum_{j=1}^k P_j \right) | \psi \rangle \\ &= \langle \psi | I | \psi \rangle = \langle \psi | \psi \rangle \\ &= \|\psi\|^2 = 1, \end{aligned}$$

$\therefore$ probabilities sum to 1.

One qubit:

$$|0\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \quad \text{bit value 0}$$

$$|1\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix} \quad \text{bit value 1}$$

CNOT on one qubit can be

$$P_0 = |0\rangle\langle 0| = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$$

"the bit is 0"

$$= \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$$

$$P_1 = |1\rangle\langle 1| = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}$$

"the bit is 1"

$$= \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}$$

$$P_0 + P_1 = I$$

measures the value of the bit.
Measuring in the computational basis.

Let P be a projector.

U unitary ($P, U \in \mathcal{L}(\mathcal{H})$)

$\Rightarrow UPU^*$ is a projector.

Proof:

$$(UPU^*)^* = U^{**} P^* U^*$$

$$= UPU^* \quad \text{and}$$

$$(UPU^*)(UPU^*)$$

$$= UPPU^* = UPU^*$$