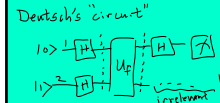
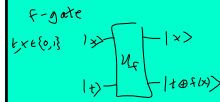


- 1. Deutsch's Problem (cont.)
- 2. Classical states, gates, circuits
- 3. Spectral Decomposition (time permitting)

1.  $f: \{0,1\} \rightarrow \{0,1\}$



Stepping through the circuit

$$|0,1\rangle \xrightarrow{H,H} \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \otimes (|0\rangle - |1\rangle)$$

$$= \frac{1}{2} (|0\rangle + |1\rangle) \otimes (|0\rangle - |1\rangle)$$

$$= \frac{1}{2} (|0\rangle \otimes (|0\rangle - |1\rangle) + |1\rangle \otimes (|0\rangle - |1\rangle))$$

$$\xrightarrow{U_f} \frac{1}{2} (|0\rangle \otimes (f(0) - f(1)) + |1\rangle \otimes (f(1) - f(0)))$$

$$\xrightarrow{H} \frac{1}{2\sqrt{2}} ((|0\rangle + |1\rangle) \otimes (f(0) - f(1)) + (|0\rangle - |1\rangle) \otimes (f(1) - f(0)))$$

$$= \frac{1}{2\sqrt{2}} (|0\rangle \otimes (f(0) - f(1)) + f(1) \otimes (|0\rangle - |1\rangle) + |1\rangle \otimes (f(0) - f(1)) - f(0) \otimes (|0\rangle - |1\rangle))$$

Case 1:  $f$  is constant:  $f(0) = f(1)$

$$|ψ\rangle = \frac{1}{2\sqrt{2}} (|0\rangle \otimes (f(0) - f(0)) + f(0) \otimes (|0\rangle - |1\rangle) + |1\rangle \otimes (f(0) - f(0)) - f(0) \otimes (|0\rangle - |1\rangle))$$

$$= \frac{1}{2\sqrt{2}} (|0\rangle \otimes (2|1\rangle) - 2|f(0)\rangle)$$

$$= |0\rangle \otimes \left( \frac{|1\rangle - |f(0)\rangle}{\sqrt{2}} \right)$$

measuring 1st qubit, get 0 with certainty.

Case 2:  $f$  is balanced,  $f(1) = \bar{f}(0)$

$$|ψ\rangle = \frac{1}{2\sqrt{2}} (|0\rangle \otimes (2|f(0)\rangle - 2|f(0)\rangle) + |1\rangle \otimes (2|f(0)\rangle - 2|f(0)\rangle))$$

measure qubit 2, get 1 with certainty.

2. An n-qubit quantum state is classical means it is a computational basis state:  $|x\rangle$  for some  $x \in \{0,1\}^n$

An n-qubit unitary operator (gate) is classical means it maps classical states to classical states.

Quantum ops that are classical must be reversible.

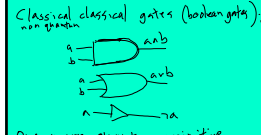
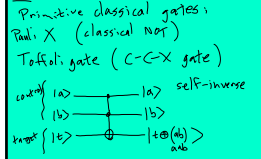
Matrix form of a classical gate is a permutation matrix:

To be classical,  $M$  must map  $e_j$  to some  $e_k$ .

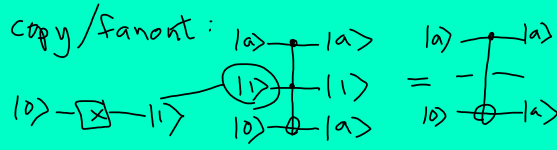
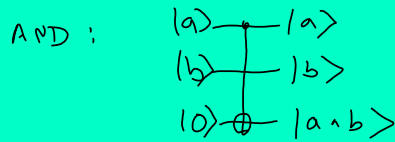
$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = \begin{bmatrix} e_2 \\ e_1 \end{bmatrix}$$

- ∴  $M$  has  $n$  many 1's exactly.
- But if 2 cols have 1 in the same row, then some row is all 0's (pigeonhole principle)
- ∴  $M$  is not invertible ( $\det = 0$ ).
- ∴ This can't happen
- ∴  $M$  has exactly one 1 in every row and every column. (def of permutation matrix)
- ∴  $M^T = M^{-1}$  - also a permutation matrix

How quantum circuits can do classical computations



One more quantum primitive: preparing a qubit in the  $|0\rangle$  state. Need to simulate AND & NOT gates quantumly.

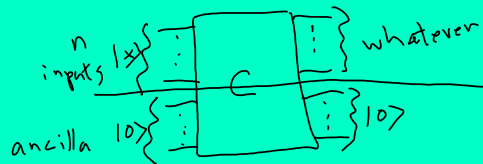


with copy, AND, NOT, can simulate any classical boolean circuit

Simulating a boolean circuit reversibly usually requires extra wires that are not input wires: ancilla qubits.

Want modularity with our quantum circuits:

Def: A quantum circuit is clean if, assuming all ancilla qubits start in the  $|0\rangle$  state, these qubits <sup>all</sup> return to the  $|0\rangle$  state at the end of the circuit



C corresponds to a unitary gate on the input qubits only.

(ancilla used only "locally" or temporarily.)

Without cleanliness, info can bleed into the ancilla, so output is entangled with the ancilla - info loss on the output.

$$f: \{0, 1\}^n \rightarrow \{0, 1\}$$

Want clean computation of  $f$ . (with just unitary gates);

